

Quadrant II - Notes

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Notes:

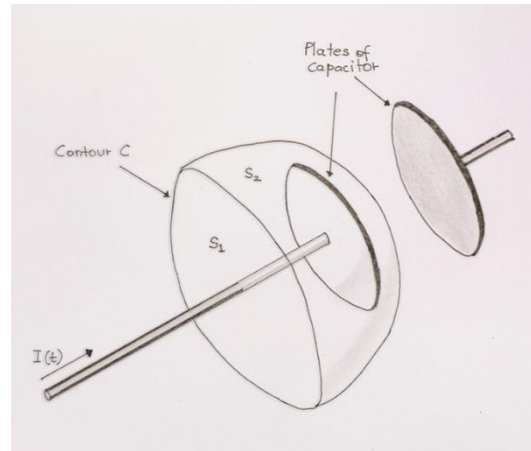
The generalization of Ampere's Law. Displacement Current

The magnetic field due to a current distribution satisfied Ampere's circuital law.

$$\oint \mathbf{H} \cdot d\mathbf{l} = \int_S \mathbf{J} \cdot \mathbf{n} d\mathbf{a} \quad - \text{Eq(1)}$$

We shall now examine this law, show that it sometimes fails, and find a generalization that is always valid.

Consider the circuit shown in figure below, which consists of a small parallel plate capacitor being charged by a constant current I



If Ampere's law is applied to the contour C and the surface S_1 , we find

$$\oint \mathbf{H} \cdot d\mathbf{l} = \int_{S_1} \mathbf{J} \cdot \mathbf{n} d\mathbf{a} = I \quad - \text{Eq(2)}$$

If on the other hand, Ampere's law is applied to the contour C and surface S_2 , then J is zero at all points on S_2 and

$$\oint \mathbf{H} \cdot d\mathbf{l} = \int_{S_2} \mathbf{J} \cdot \mathbf{n} d\mathbf{a} = 0 \quad - \text{Eq(3)}$$

Equations (2) and (3) contradict each other and thus cannot both be correct. If C is imagined to be a great distance from the capacitor, it is clear that the situation is not substantially different from the standard Ampere's law cases. One is thus led to think that Eq(2) is correct, since it is not dependent of the new feature, namely the capacitor. Eq(3), on the other hand, requires consideration of the capacitor for its deduction. It would appear, then, that Eq(3) requires modification.

$$\nabla \times \mathbf{H} = \mathbf{J} \quad - \text{Eq(4)}$$

this also requires modification.

The proper modification can be made by noting that Eq(2) and Eq(3) give different results because the integrals on the right-hand sides are different.

Phrased mathematically,

$$\int_{S_2} \mathbf{J} \cdot \mathbf{n}_2 d\mathbf{a} - \int_{S_1} \mathbf{J} \cdot \mathbf{n}_1 d\mathbf{a} \neq 0 \quad - \text{Eq(5)}$$

S_1 and S_2 together form a closed surface (they join at C); however, n_2 is outward drawn and n_1 inward drawn. If this fact is taken into account, Eq(5) may be written

$$\oint_{S_1+S_2} \mathbf{J} \cdot \mathbf{n} \, d\mathbf{a} \neq \mathbf{0} \quad - \text{Eq(6)}$$

Phrased physically, the net transport current through the closed surface $S_1 + S_2$ does not vanish because charge is piling up on the plate of the condenser enclosed by the surface. Charge conservation requires, according to (ENTER EQUATION HERE),

$$\oint_{S_1+S_2} \mathbf{J} \cdot \mathbf{n} \, d\mathbf{a} = - \int_V \frac{\partial \rho}{\partial t} \, dv \quad - \text{Eq(7)}$$

Because inside the volume V enclosed by $S_1 + S_2$ the charge density ρ is changing with time on the condenser plate. In differential form Eq(7) is expressed by the equation of continuity,

$$\nabla \cdot \mathbf{J} + \frac{\partial \rho}{\partial t} = \mathbf{0} \quad - \text{Eq(8)}$$

It is now clear what is wrong with Eq(4): Taking its divergence we have

$$\nabla \cdot \nabla \times \mathbf{H} = \mathbf{0} = \nabla \cdot \mathbf{J},$$

since the divergence of a curl is identically zero. Thus the relation $\nabla \cdot \mathbf{J} = \mathbf{0}$, as implied by Eq(4), is consistent with charge conservation in the present situation, and so something must be added to the right side of Eq(4) that will give $\partial \rho / \partial t$ in Eq(8). What this could be is seen from the relation of ρ to electric displacement:

$$\nabla \cdot \mathbf{D} = \rho \quad - \text{Eq(9)}$$

Inserting ρ from the Eq(9) into Eq(8), we get

$$\nabla \cdot \mathbf{J} + \frac{\partial}{\partial t} \nabla \cdot \mathbf{D} = \nabla \cdot \left(\mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \right) = \mathbf{0}.$$

If $\partial \mathbf{D} / \partial t$ is added to Eq(4), then its divergence will correctly give Eq(8). (ENTER STATEMENT HERE). Inclusion of $\partial \mathbf{D} / \partial t$ gives the generalized Ampere's law:

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t}. \quad - \text{Eq(10)}$$

The introduction of the second term on the right, which is known as the displacement current, represents one of Maxwell's major contributions to electromagnetic theory.
